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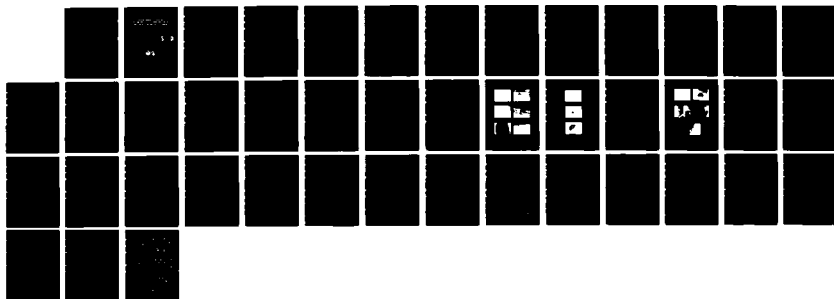
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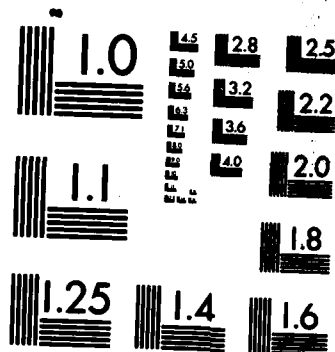
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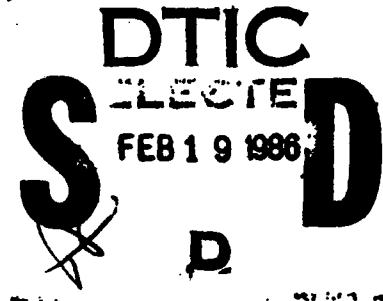
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Characterization of Aluminum Nitride/Aluminum Oxide Reactively Sputtered Antireflection Coatings

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SEPTEMBER 1985



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FOREWORD

The work described in this report was performed by the Naval Weapons Center, China Lake, Calif., under Arpa Order 3343. This report is a preliminary study of the multilayer system $\text{AlN}/\text{Al}_2\text{O}_3$ and its resistance to certain environmental conditions. The work described was carried out during FY 1984.

The report has been reviewed for technical accuracy by A. D. Baer and J. L. Stanford.

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19. ABSTRACT (Continue on reverse if necessary and identify by block number) <p>(U) Aluminum nitride/aluminum oxide multilayer protective antireflection coatings are being developed for use on laser windows. These materials have been shown to be relatively stable and scratch resistant, thus demonstrating their applicability for this use. Both two- and three-layer coatings were designed at the wavelength of 0.500μm. The design techniques utilized, as well as the resulting transmission spectra, are discussed. Single-layer films have been deposited on fused quartz and calcium fluoride substrates, and a three-layer film has been deposited on a fused quartz substrate. Studies of the resistance of these films to various environmental conditions including laser radiation, fluorine, electron beam irradiation, and humidity are presented. Surface analysis techniques such as scanning Auger microscopy (SAM), scanning electron microscopy (SEM), energy dispersive X-ray microanalysis (EDX), X-ray photoelectron spectroscopy (XPS), and Nomarski microscopy were used to examine coating composition and stoichiometry. Preliminary measurements on the laser resistance of the materials indicate that damage at defects is the predominant failure mode.</p>					
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CONTENTS

Introduction.....	3
Multilayer Designs.....	4
Two-Layer Design.....	4
Three-Layer Design.....	7
Sample Preparation.....	11
Film Structure.....	16
Results.....	19
Laser Radiation.....	19
Fluorine Exposure.....	24
Electron Beam Irradiation.....	24
Humidity.....	24
Conclusions.....	26
References.....	27
Appendix:	
Computer Program ANTIREF for Calculating Three-layer.....	31
Antireflection Coating Solutions.	

Figures:

1. Two-layer Design Configuration.....	6
2. Theoretical Transmittance of Two-layer Design.....	6
3. Normalized Electric Field Intensity of Two-layer Design.....	6
4a. Antireflection Solutions for a Three-layer Design with Fused Quartz as Substrate.....	9
4b. Antireflection Solutions for a Three-layer Design with Calcium Fluoride as Substrate.....	9
5. Three-layer Design Configuration with Fused Quartz as Substrate.....	9
6. Theoretical Transmittance of Three-layer Design with Fused Quartz as Substrate.....	9
7. Normalized Electric Field Intensity of Three-layer Design with Fused Quartz as Substrate.....	10
8. Experimental Transmittance of Three-layer Design with Fused Quartz as Substrate.....	15
9a. Topography of AlN Single Layer.....	17
9b. Cross Section of AlN Single Layer.....	17



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10a. Topography of Al_2O_3 Single Layer.....	17
10b. Cross Section of Al_2O_3 Single Layer.....	17
11a. Topography of $\text{AlN}/\text{Al}_2\text{O}_3/\text{AlN}$ Multilayer.....	17
11b. Cross Section of $\text{AlN}/\text{Al}_2\text{O}_3/\text{AlN}$ Multilayer.....	17
12a. Enlargened Top View of AlN Single Layer Defect.....	18
12b. Enlargened Top View of Al_2O_3 Single Layer Defect.....	18
12c. Enlargened Top View of $\text{AlN}/\text{Al}_2\text{O}_3/\text{AlN}$ Multilayer Defect.....	18
13a. Damage Site 1 Showing Small Damage Craters Amongst Defects.....	20
13b. Damage Site 2 Showing More and Larger Damage Craters Amongst Defects (than Site 1).....	20
13c. Damage Site 3 Showing a Reduced Number of Apparent Defects and More Damage Craters (than Site 2).....	20
13d. Damage Site 4 Showing Agglomeration of Damage Craters.....	20
13e. Damage Site 5 Showing Small Pitmarks and Loss of Much of Outer Film.....	20
14a. AES Spectrum of Smooth Film Portion of Site 2.....	22
14b. AES Spectrum Taken Within Damage Crater of Site 2.....	22
15. AES Spectrum of Smooth Film Portion of Site 4.....	23
16a. AES Profile of AlN Single Layer Exposed to Fluorine.....	25
16b. AES Profile of Al_2O_3 Single Layer Exposed to Fluorine.....	25

Tables:

1. Information on Samples Prepared in dc Magnetron System.....	12
2. Information on Samples Prepared in rf Diode System.....	13
3. Purities of Materials Used in Sample Depositions.....	13
4. Laser Energy Densities for the Five Damage Sites.....	21

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INTRODUCTION

New multilayer antireflective coatings are being developed for protective use on laser windows. These coatings consist of alternating layers of aluminum nitride and aluminum oxide, which are stable and scratch resistant. Both materials have large bandgaps and are highly transparent from the ultraviolet to the infrared, thus making them ideal for use with many laser systems.

Both two- and three-layer designs have been developed for use at certain wavelengths. The two-layer design was based upon the work of Cox and Hass (Reference 1) and uses aluminum oxide as the outer material. This is particularly desirable since Al_2O_3 is resistant to moisture as well as to fluorine. The three-layer design was developed by applying the works of Baer and Mouchart who have determined techniques which allow for much design flexibility (References 2 and 3, respectively). This flexibility is the result of an extra degree of freedom which may be used to optimize the design with respect to low absorption, large bandwidth, good adherence, or ease of fabrication (Reference 4).

Preparations of the films have been achieved by numerous techniques including reactive sputtering and evaporation, chemical vapor deposition, ion plating, activated decomposition of organometallic compound vapors, aluminum anodization, and laser-assisted deposition during electron beam evaporation (References 5 through 36). In this study, the methods of dc magnetron, ion beam, and rf diode sputtering have been used to deposit single layers and multilayers. Films prepared by these techniques have typically been shown to be stoichiometric and of high packing density.

Studies on the reactions of the multilayer materials with fluorine, electron beam irradiation, and humidity as well as on their damage by laser radiation will be discussed. The main techniques used for analysis and characterization of the films after exposure to the above environments are scanning Auger microscopy (SAM), scanning electron microscopy (SEM), energy dispersive X-ray microanalysis (EDX), X-ray photoelectron spectroscopy (XPS), and Nomarski microscopy.

MULTILAYER DESIGNS

The choice of materials, AlN and Al₂O₃, necessitate the use of multiple-layer design techniques for obtaining zero reflectance. Both materials have high transmittances at the wavelength of interest (i.e., 0.500μm), but by themselves cannot accomplish the reflectance requirement. These materials have been prepared with indices of refraction which allow for production of zero reflectance coatings.

Since quarter-wave designs which produce zero reflectance are not possible, general solutions to the two- and three-layer cases must be used. These solutions allow flexibility in the choices of layer thicknesses, which is very desirable from the standpoints of absorption in the layers and film preparation.

Absorption in the layers is a problem which normally must be accounted for when designing multilayer systems, but the design techniques utilized in this study do not include this factor. Thus, an indirect approach should be used such as optimizing the solution through manipulations of layer thicknesses (Reference 2).

TWO-LAYER DESIGN

As discussed by Cox and Hass (Reference 1), for zero reflectance, when the multilayer consists of two layers whose thicknesses are not related by integral multiples of each other, one obtains:

$$\frac{\tan \phi_1}{\tan \phi_2} = - \frac{r_1 - r_3 + r_2(1 - r_1 r_3)}{r_1 - r_3 - r_2(1 - r_1 r_3)} \quad (1)$$

$$\text{and} \quad \tan \phi_1 \tan \phi_2 = \frac{r_1 + r_3 + r_2(1 + r_1 r_3)}{r_1 + r_3 - r_2(1 + r_1 r_3)} \quad (2)$$

$$\text{where} \quad \phi_j = \frac{2\pi n_j l_j}{\lambda} \equiv 2\pi X_j ,$$

$$\text{and} \quad r_j = \frac{n_{j-1} - n_j}{n_{j-1} + n_j} .$$

Note that n_j is the index of refraction of the j th layer, l_j is its thickness, and λ is the design wavelength.

Equations (1) and (2) may be written in the form:

$$X_1 = \frac{1}{2\pi} \tan^{-1} \left[\pm \left(\frac{n_1^2(n_3-n_0)(n_2^2-n_0n_3)}{(n_1^2n_3-n_2^2n_0)(n_0n_3-n_1^2)} \right)^{1/2} \right] + \frac{m_1}{2} \quad (3)$$

$$\text{and } X_2 = \frac{1}{2\pi} \tan^{-1} \left[\pm \left(\frac{n_2^2(n_3-n_0)(n_0n_3-n_1^2)}{(n_1^2n_3-n_2^2n_0)(n_2^2-n_0n_3)} \right)^{1/2} \right] + \frac{m_2}{2}, \quad (4)$$

where $m_1 = 0, \pm 1, \pm 2, \dots$,

and $m_2 = 0, \pm 1, \pm 2, \dots$.

Thus, one may obtain a series of X_1 's and X_2 's and hence a series of physical thicknesses, l_1 's and l_2 's. Note that equations (3) and (4) imply four possible pairs of solutions for each m_1 and m_2 set, but in actuality only two of these pairs will satisfy equations (1) and (2).

Applying equations (3) and (4) to the present case of interest, where $n_0=1.00$, $n_1=1.60$, $n_2=2.00$, $n_3=1.46$ (these are the incident medium, outer aluminum oxide layer, inner aluminum nitride layer, and fused quartz substrate indices of refraction, respectively), and $\lambda=0.500\mu\text{m}$, a zero reflectance choice of thicknesses is then given by $l_1=0.0631\mu\text{m}$ and $l_2=0.0832\mu\text{m}$. This configuration is shown in Figure 1. From this information, and assuming no absorption in the layers, normal light incidence, and constant indices of refraction for the wavelength region 0.3 to $0.9\mu\text{m}$, the transmission spectrum has been determined using a program developed by Loomis (Reference 37).

This spectrum is shown in Figure 2 as a function of wavelength. As can be seen, there is 100% transmission at $0.500\mu\text{m}$ and a bandwidth of approximately $0.06\mu\text{m}$ with a reflectance loss of 0.5% or less. The normalized electric field intensity as a function of depth into the layers is shown in Figure 3 and was determined by Elson (Reference 38). It is apparent that the electric field minimum occurs in the thicker aluminum nitride layer. This is advantageous from the standpoint that the aluminum nitride layer is the more absorbing material as shown by preliminary absorption measurements.

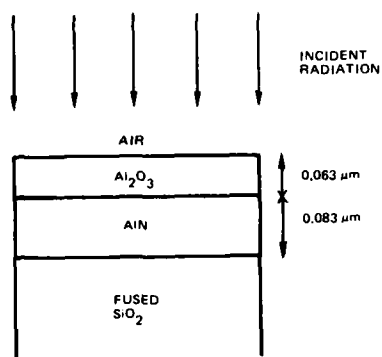


FIGURE 1. Two-layer Design Configuration.

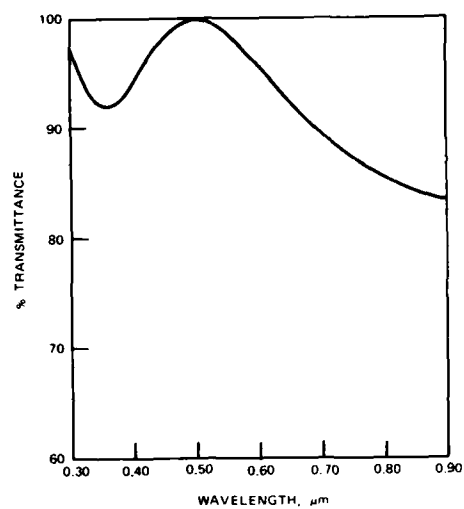


FIGURE 2. Theoretical Transmittance of Two-layer Design.

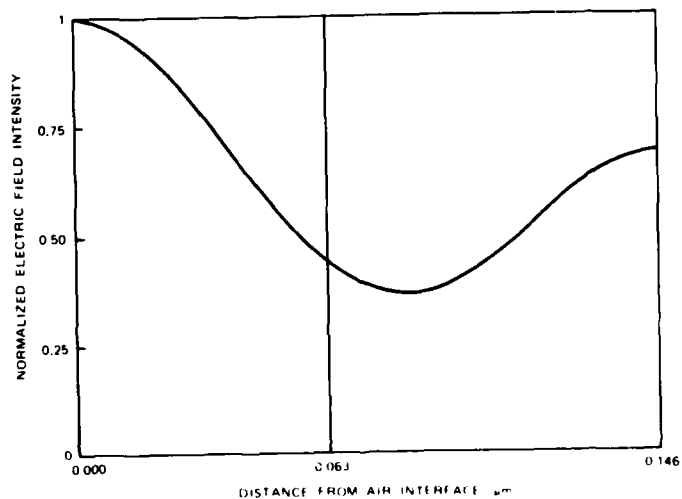


FIGURE 3. Normalized Electric Field Intensity of Two-layer Design.

THREE-LAYER DESIGN

Based upon the works of Baer and Mouchart (References 2 and 3, respectively), an analytic solution for a general three-layer antireflective coating has been utilized. According to Baer, Smith's theorem may be used when determining what conditions must be applied to obtain zero reflectance. These conditions indicate that both the real and imaginary components of the Fresnel coefficients must vanish. Then, as discussed by Baer, one obtains:

$$X_1 = \pm \frac{1}{4\pi} \cos^{-1} \left[\frac{r_2^2 + r_1^2 - |r_b|^2(1 + r_2^2 r_1^2)}{2r_2 r_1 (|r_b|^2 - 1)} \right] + \frac{m_1}{2}, \quad (5)$$

and
$$X_2 = \frac{1}{4\pi} (\delta_a + \delta_b) + \frac{m_2}{2}, \quad (6)$$

where
$$r_a = \frac{r_2 + r_1 \exp(-4\pi X_1 i)}{1 + r_2 r_1 \exp(-4\pi X_1 i)}, \quad (7)$$

$$r_b = \frac{r_3 + r_4 \exp(-4\pi X_3 i)}{1 + r_3 r_4 \exp(-4\pi X_3 i)}, \quad (8)$$

and
$$X_j = \frac{n_j l_j}{\lambda}, \quad (9)$$

$$m_1 = 0, \pm 1, \pm 2, \dots,$$

$$m_2 = 0, \pm 1, \pm 2, \dots$$

Note that δ_a and δ_b refer to the phases of r_a and r_b respectively and that the X_j and r_j are as defined earlier. Hence, for a choice of X_3 , and given the indices of refraction of the layers, one may obtain a series of solutions. Note that there are two series of solutions for each X_3 .

A computer program in the language of Fortran has been written which calculates and then plots the solutions of X_1 and X_2 as functions of X_3 given the indices of refraction and design wavelength. A listing of this program may be found in the Appendix.

Figure 4a shows a plot of X_1 and X_2 as functions of X_3 for $\lambda=0.500\mu\text{m}$, $n_0=1.00$, $n_1=2.00$, $n_2=1.60$, $n_3=2.00$, and $n_4=1.46$. Figure 4b shows a similar plot with $\lambda=0.500\mu\text{m}$, $n_0=1.00$, $n_1=2.00$, $n_2=1.60$, $n_3=2.00$, and $n_4=1.44$. These indices of refraction correspond to the incident medium, aluminum nitride layer, aluminum oxide layer, aluminum nitride layer, and fused quartz substrate (Figure 4a) or calcium fluoride substrate (Figure 4b), respectively. Note that the solid curves correspond to the first solution and the dashed to the second.

A particular solution from the results shown in Figure 4a is $X_3=0.250$, $X_1=0.447$, and $X_2=0.326$. Hence, $l_1=0.112\mu\text{m}$, $l_2=0.102\mu\text{m}$, and $l_3=0.625\mu\text{m}$. This configuration is shown in Figure 5.

As in the two-layer case, a plot of the transmission spectrum as a function of wavelength is shown in Figure 6 (for the particular solution discussed above). Again, there is 100% transmission at $0.500\mu\text{m}$ as expected, but the bandwidth at a reflectance loss of 0.5% is reduced to approximately $0.02\mu\text{m}$. Hence, from the standpoint of bandwidth, this design is not as desirable. The normalized electric field intensity is shown in Figure 7. In this case, the electric field minimum again occurs in one of the aluminum nitride layers, with the maximum occurring in the aluminum oxide layer. Thus, as in the two-layer case, the design appears to be a favorable choice from the standpoint of absorption.

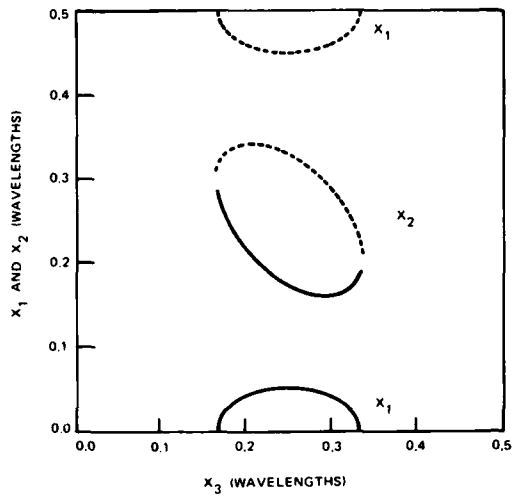


FIGURE 4. Antireflection Solutions for a Three-layer Design with Fused Quartz as Substrate. (a)

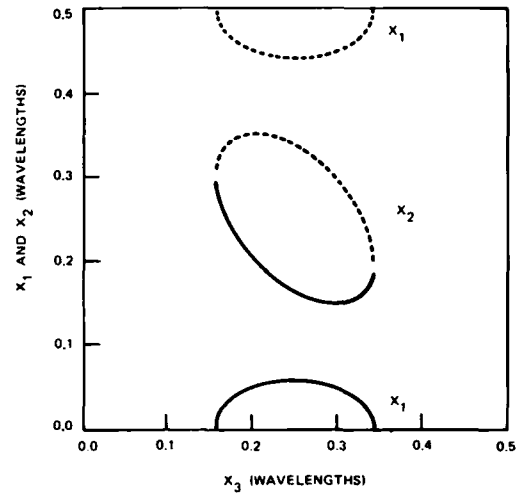


FIGURE 4. Antireflection Solutions for a Three-layer Design with Calcium Fluoride as Substrate. (b)

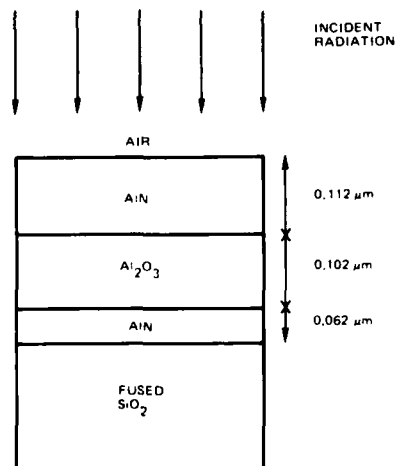


FIGURE 5. Three-layer Design Configuration with Fused Quartz as Substrate.

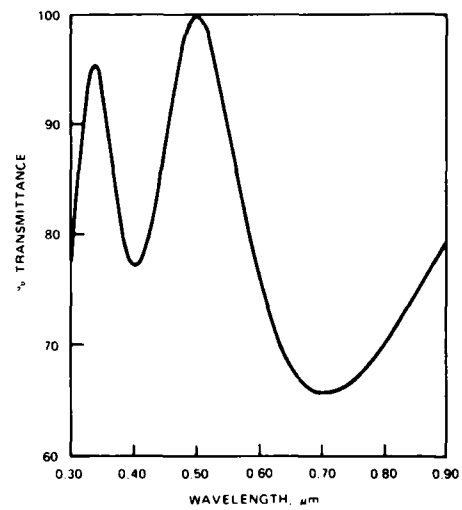


FIGURE 6. Theoretical Transmittance of Three-layer Design with Fused Quartz as Substrate.

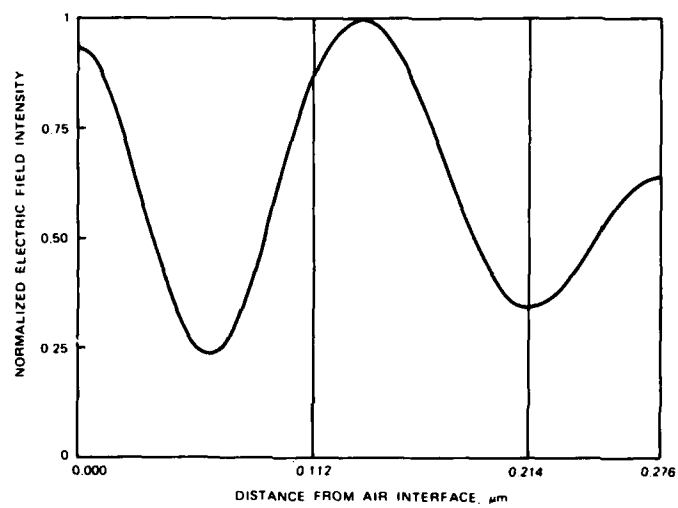


FIGURE 7. Normalized Electric Field Intensity of Three-layer Design with Fused Quartz as Substrate.

SAMPLE PREPARATION

Several single-layer samples and one three-layer sample were deposited onto substrates of fused quartz, calcium fluoride, or silicon. Substrate diameters were 1 inch and thicknesses were 0.063 inch, 0.058 inch, and 0.023 inch, respectively. Very thin layers of Ag/Cr/Mo, Au/Cr/Mo, or Au/Cr were deposited onto some of the substrates prior to AlN or Al₂O₃ deposition for the purpose of aiding in eliminating charging problems during analysis. Tables 1 and 2 contain lists of the specimens prepared along with any pertinent information relating to their depositions for depositions in the dc magnetron and rf diode sputtering systems, respectively. Approximate deposition temperatures for the two systems were 160°C and 300°C, respectively. Note that samples on Si or on aiding layers were used as witnesses. Sample M-177 of Al₂O₃, prepared in the ion beam sputtering system, was deposited on aiding layers of Au/Cr with a substrate of Si. The approximate gas pressure was 1.5×10^{-4} torr of Ar, deposition was at room temperature, and the estimated film thickness was 2800Å. Table 3 contains information regarding the purities of the materials used (for the various depositions) and their suppliers.

Most of the samples were deposited using a dc magnetron sputtering system with a 5-inch Varian sputter gun (samples listed with a "T" before their numbers; see Table 1). Some samples were prepared in a Randex model 2400 rf diode sputtering system (samples with an "S" before their numbers; see Table 2). One sample was prepared by ion beam deposition using an Ion Technology, Inc., 2.5mm ion gun (as indicated by an "M" before the sample number above).

TABLE 1. Information on Samples Prepared in dc Magnetron System.

Sample Number	Film Material	Aiding Layers	Approximate Gas Pressure (torr)	Substrate	Estimated Film Thickness (Å)
T-261S	AlN	-	$3 \times 10^{-3} \text{N}_2$	Fused Quartz	2030
T-265S	AlN	-	$1.5 \times 10^{-3} \text{N}_2$ $1.5 \times 10^{-3} \text{Ar}$	Si	2000
T-266S	AlN	Ag/Cr/Mo	$3 \times 10^{-3} \text{N}_2$	Fused Quartz	2000
T-268S	AlN	Ag/Cr/Mo	$1.5 \times 10^{-3} \text{N}_2$ $1.5 \times 10^{-3} \text{Ar}$	Fused Quartz	2000
T-270S	Al ₂ O ₃	-	$3 \times 10^{-4} \text{O}_2$ $2.7 \times 10^{-3} \text{N}_2$	Fused Quartz	2000
T-276S	AlN	-	$1.5 \times 10^{-3} \text{N}_2$ $1.5 \times 10^{-3} \text{Ar}$	CaF ₂	2000
T-281S	AlN	-	$3 \times 10^{-3} \text{N}_2$	CaF ₂	4000
T-289S	(a) AlN (outer)	-	(a) $3 \times 10^{-3} \text{N}_2$	Fused Quartz	(a) 1120
	(b) Al ₂ O ₃ (middle)	-	(b) $3 \times 10^{-4} \text{O}_2$ $2.7 \times 10^{-3} \text{N}_2$		(b) 1020
	(c) AlN (inside)	-	(c) $3 \times 10^{-3} \text{N}_2$		(c) 625
T-290S	AlN	Ag/Cr/Mo	$3 \times 10^{-3} \text{N}_2$	Fused Quartz	2000
T-290S	AlN	Au/Cr/Mo	$3 \times 10^{-3} \text{N}_2$	Fused Quartz	2000
T-297S	AlN	Au/Cr	$3 \times 10^{-3} \text{N}_2$	Si	2800

TABLE 2. Information on Samples Prepared in rf Diode System.

Sample Number	Film Material	Approximate Gas Pressure (torr)	Substrate	Film Thickness (Å)
S-826	Al ₂ O ₃	8×10^{-4} O ₂ 7.2×10^{-3} N ₂	Fused Quartz	2100
S-827	AlN	2×10^{-3} N ₂ 6×10^{-3} Ar	Fused Quartz	7000
S-833	Al ₂ O ₃	8×10^{-4} O ₂ 7.2×10^{-3} Ar	Fused Quartz	2100
S-834	AlN	2×10^{-3} N ₂ 6×10^{-3} Ar	Fused Quartz	7000
S-835	AlN	8×10^{-4} N ₂ 7.2×10^{-3} Ar	Fused Quartz	7000

TABLE 3. Purities of Materials Used in Sample Depositions.

Material	Purity (%)	Source
Al	99.999	Specialty Metals Division of Varian
Al ₂ O ₃	Unknown	Microelectronics Branch, NAVWPNCEN
Au	99.999	American Smelting Company
Mo	99.999	American Smelting Company
Cr	99.996	Specialty Metals Division of Varian

The three-layer sample was prepared according to the 0.500 μ m design discussed previously (particular solution). Indices were assumed to be 2.00 and 1.60 for AlN and Al₂O₃, respectively, and the materials were deposited until the necessary thicknesses had been obtained. Although the indices of refraction have not been accurately determined in this study, many researchers have previously determined them to range from 1.70-2.15 for AlN and 1.42-1.79 for Al₂O₃ (References 5, 6, 16, 23, 25, 26, 28, 29, 30, 35, and 36), implying that the assumed values are realistic. Furthermore, recent laboratory studies indicate that the assumed indices are more likely.

A transmission spectrum of the three-layer sample coating over the wavelength region of 0.3 to 0.9 μ m was measured using a Beckman DU-7 UV/VIS spectrophotometer. This spectrum is shown in Figure 8. In comparison to the theoretical spectrum shown in Figure 6, it is apparent that the two are similar in curve shape, but some additional tuning is necessary in the experimental sample preparation to more closely approximate the theoretical curve. The experimental maximum in transmission of 98.8% occurred at 0.553 μ m. By the use of a program written by Loomis (Reference 37), variations in the thicknesses and indices of the layers were made to attempt to determine what changes would cause a calculated transmission spectrum to approach the appearance of the experimental spectrum. It was found that the differences between the two spectra is a combination of a need to determine the actual film indices and to measure more accurately the deposition rates of the materials. At this time, efforts are being made to accomplish these tasks.

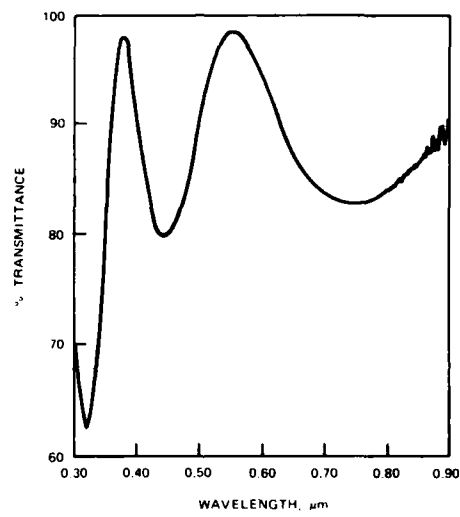


FIGURE 8. Experimental Transmittance of Three-layer Design with Fused Quartz as Substrate.

FILM STRUCTURE

Study of the structures of both single and multilayer films were made using a scanning electron microscope. The films were coated with thin layers of Au (approximately 150Å) to avoid charging problems.

Shown in Figures 9a and 9b are typical top and cross-sectional views of an AlN single layer. This film was sample number T-297S prepared in the dc magnetron system. Figures 10a and 10b show top and cross-sectional views of an Al₂O₃ single layer prepared in the ion beam system (sample M-177). Figures 11a and 11b show top and cross-sectional views of the AlN/Al₂O₃/AlN multilayer (sample T-289S) prepared in the dc magnetron system. In the last case, the three layers can be distinguished.

Many defects are apparent on the surface of the AlN single and multilayer films. It is not known at this time whether the defects are growth defects of a nodular type or whether they are bubbles or splatters on the surfaces only. These defects are typically on the order of 5-10µm in diameter. The Al₂O₃ specimen is of a much finer surface structure, although smaller defects (approximately 1µm in diameter) are apparent on it also. Figures 12a through 12c show magnified views of the AlN, Al₂O₃, and multilayer defects, respectively.

The cross sections or side views of the AlN layers display columnar growth type structures. It has been determined by transmission electron diffraction that these films are polycrystalline. The side views of the Al₂O₃ layers show much smoother textures, with no columnar growth structures apparent at the magnifications achieved in these experiments. A pattern was not obtained from the diffraction studies for the Al₂O₃, and hence it is probably amorphous. Note that the Al₂O₃ single layer and multilayer specimens were prepared in different sputtering systems, but the texture of the Al₂O₃ still remains similar in the two cases.

From electron spectroscopy for chemical analysis (ESCA) studies, AlN materials have been determined to be 95% stoichiometric with the remainder consisting mostly of Al₂O₃. Al₂O₃ films have been shown to be 99% stoichiometric.

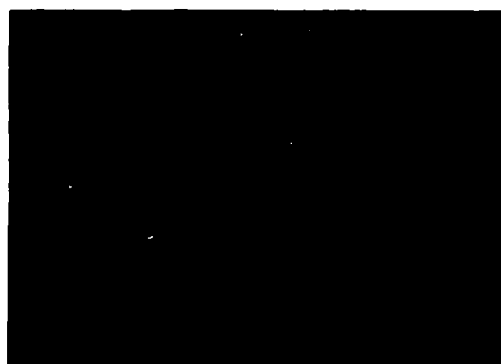


FIGURE 9. Topography of AlN Single Layer.
(a)



FIGURE 9. Cross Section of AlN Single Layer.
(b)



FIGURE 10. Topography of Al₂O₃ Single Layer.
(a)



FIGURE 10. Cross Section of Al₂O₃ Single Layer.
(b)

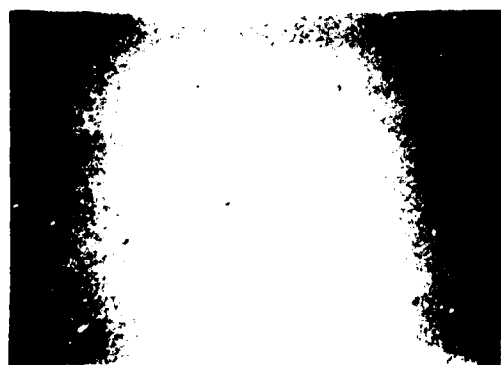


FIGURE 11. Topography of AlN/Al₂O₃/AlN Multilayer.
(a)



FIGURE 11. Cross Section of AlN/Al₂O₃/AlN Multilayer.
(b)



FIGURE 12. Enlarged Top View of AlN Single Layer Defect.
(a)

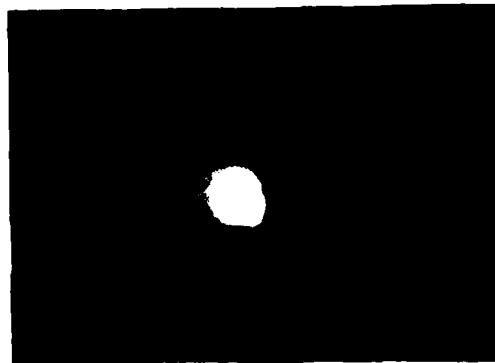


FIGURE 12. Enlarged Top View of Al₂O₃ Single Layer Defect.
(b)



FIGURE 12. Enlarged Top View of AlN/Al₂O₃/AlN Multilayer Defect.
(c)

RESULTS

Results have been obtained on sample behavior following exposure to laser radiation, fluorine, electron beam irradiation, and humidity. The samples studied were discussed earlier, and will be referred to by their designated sample numbers. Not all samples given earlier will be discussed since some were only used for obtaining preliminary spectral information and since experimental difficulties were encountered during analysis with others which were associated with film quality and charging problems.

LASER RADIATION

The effects of laser radiation on the multilayer sample were studied with a Phi 600 scanning Auger microscopy system (after damage). The sample surface was cleaned during analysis by sputtering with Ar ions using an accelerating voltage of 4kV. For SiO_2 , an approximate removal rate has been determined to be 40Å/min.

A pulsed dye laser in a triaxial configuration was utilized for damaging the film at a wavelength of $0.497\mu\text{m}$. The first step in the laser damage study was to scratch crosshairs onto the film on opposite sides of its center along a sample diagonal. These crosshairs were then used as index marks for locating five equally spaced sites between them. Each site was then irradiated in air with a few laser pulses of increasing energy density. The pulse shape was a 1mm diameter flat top ($\pm 5\%$ flatness) with a duration of $0.55\mu\text{s}$. Film breakdown thresholds were roughly determined by observing the beginning of film changes with a video microscopy system (VIMS). Note that no attempts were made in this study to determine actual damage thresholds of the materials. The main purpose of this study was to determine the behavior of the film after exposure to laser radiation.

Table 4 lists the energy density of each shot for the five sites. Shown in Figures 13a through 13e are portions of the centers of each of the five sites. These figures in comparison to a typical undamaged site (such as that in Figure 11a) show an increasing number of damage craters with respect to defects for an increasing energy density. Site 1 shows mostly defects plus some small damage craters. Sites 2 and 3 show more and larger damage craters with a reduced number of apparent defects. Site 4 shows the beginning of damage crater agglomeration, and Site 5 shows many small pit marks which appear to lay at the centers of what were originally damage craters (the pits are at the centers of circular patterns which resemble agglomerated craters). In addition, for Site 5, it is apparent that there is a loss of much of the outer film.



FIGURE 13. Damage Site 1 Showing Small Damage Craters Amongst Defects.
(a)



FIGURE 13. Damage Site 2 Showing More and Larger Damage Craters Amongst Defects (than Site 1).
(b)



FIGURE 13. Damage Site 3 Showing a Reduced Number of Apparent Defects and More Damage Craters (than Site 2).
(c)

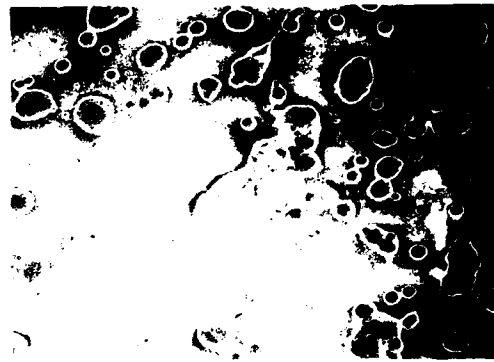


FIGURE 13. Damage Site 4 Showing Agglomeration of Damage Craters.
(d)



FIGURE 13. Damage Site 5 Showing Small Pitmarks and Loss of Much of Outer Film.
(e)

TABLE 4. Laser Energy Densities for the Five Damage Sites.

Energy Density (J/cm ²)					
	Site 1	Site 2	Site 3	Site 4	Site 5
1st pulse	3.2	9.2	12.	20.	50.
2nd pulse	5.9	9.2	-	-	-
3rd pulse	9.4	-	-	-	-

The approximate number of defects for typical portions of the film where no damage was done is 70 (see the sampling area of Figure 11a for an example). For Sites 1 and 2, the number of damage craters plus defects which did not damage are approximately 60 and 70 (in the same size sampling area), respectively. These are similar to that of the undamaged site, which implies that damage is initiating at defects. For Sites 3 through 5, the number of craters plus undamaged defects is large in comparison to the number of defects on the undamaged site, increasing in number with increasing energy density (assuming the pit marks of Site 5 are indeed the centers of what were originally damage craters), although few apparent defects remain on the sites. These results support damage initiating at the defect sites, with some other possible mode(s) of damage (e.g., caused by non-observable defects) also contributing to the cratering.

Auger spectra were obtained for some of the damage sites. Thin layers of gold were deposited on the films prior to the study and then removed by Ar ion sputtering (as discussed earlier) just until Auger peaks corresponding to the underlying film were obtained. Figures 14a and 14b show spectra obtained at Site 2. Figure 14a illustrates a portion of the damage site which appeared to be unaffected by the laser radiation (herein referred to as smooth film) and Figure 14b was taken within a damage crater. It is apparent that the smooth film consists mostly of AlN and the crater site consists of AlN plus Al₂O₃. Similar spectra were obtained for some of the other sites, where the spectra from the craters always showed an increase in Al₂O₃ with respect to AlN. In addition, when comparing damage sites, it is apparent that the amount of Al₂O₃ with respect to AlN increased with increasing energy density. Figure 15 shows the spectrum obtained on a smooth portion of the film at Site 4. A possible explanation for the growth of the oxide peak is the dissociation of some of the surface AlN, and then combination of the Al with oxygen to form Al₂O₃ (recall that these studies were performed in air). Thus, the amount of oxidation of AlN would increase with energy density. In addition, the larger oxide peaks in the craters may be explained by assuming that these are positions on the film where the energy was absorbed more, or not conducted away as quickly, hence resulting in temperature buildup and more oxidation.

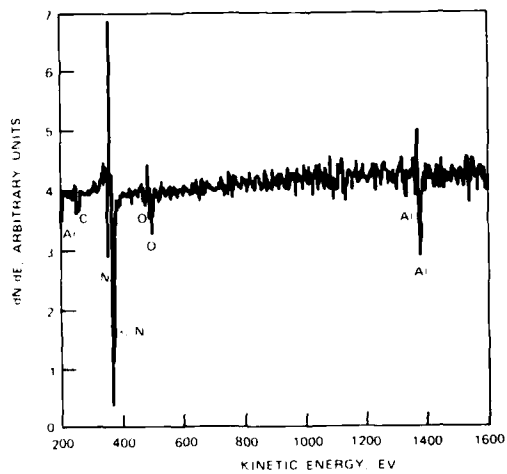


FIGURE 14. AES Spectrum of Smooth Film Portion of Site 2.
(a)

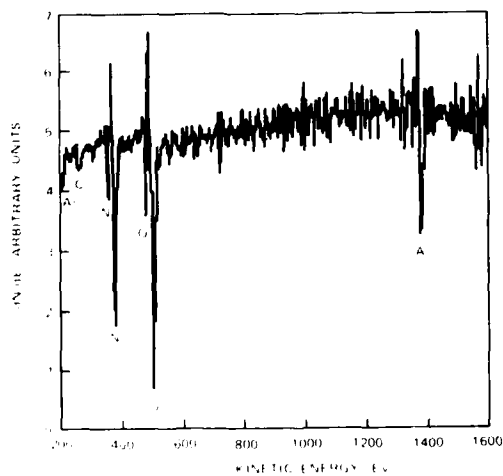


FIGURE 14. AES Spectrum Taken Within Damage Crater of Site 2.
(b)

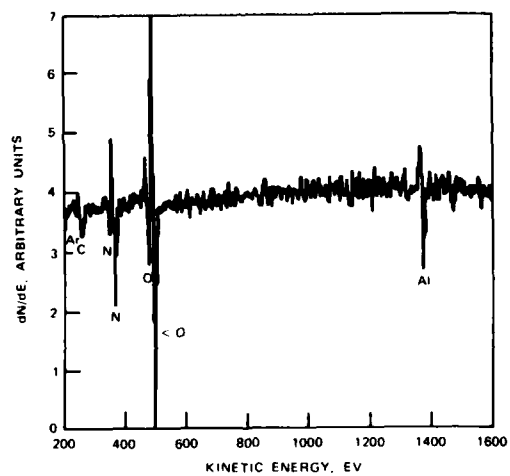


FIGURE 15. AES Spectrum of Smooth Film Portion of Site 4.

FLUORINE EXPOSURE

Single layer samples T-297S and M-177 of AlN and Al₂O₃, respectively, were exposed to 10% HF in Ar for 1 hour, and were then examined for fluorine penetration.

EDX was initially used to detect the fluorine. Since this technique examines the constituents of the materials to a much greater depth (on the order of a micron) than a technique such as AES, any severe fluorine penetration should be detected. Fluorine was not detected in either the AlN or Al₂O₃. Hence, either the fluorine had not penetrated the material to a very great depth, or the fluorine concentration was low (low atomic number materials have low X-ray intensities due to low fluorescence yield and high sample absorption; see Reference 39).

AES and XPS were then used to analyze the film makeup since the results of the EDX had indicated no great depth penetration of the fluorine in either film. With these techniques, fluorine was observed. Figures 16a and 16b display AES profiles of the AlN and Al₂O₃ films, respectively. Note that the AES system used in this case is a Physical Electronics 548 Auger/ESCA system, which employed a 1kV beam of Ar ions for sputter removal. For SiO₂, a removal rate of 15Å/min. is typical. Assuming this rate for AlN and Al₂O₃ yields approximate depth penetrations of 30-45Å for both films, although the fluorine appears to drop off (with depth) at a slightly higher rate in the Al₂O₃ film.

ELECTRON BEAM IRRADIATION

During surface studies such as SEM and SAM, physical changes of the AlN and Al₂O₃ films were observed. Dark spots indicating changes in electron yields were observed at positions on the films where the electron beams were concentrated. At this time it is uncertain whether these changes were temporary (i.e., caused by contaminants being deposited on or removed from the film surfaces) or permanent. It is important to note that, as discussed earlier, it was necessary to coat the films with thin layers of a conducting material such as Au to avoid charging problems.

HUMIDITY

Single layer samples T-297S and M-177 of AlN and Al₂O₃, respectively (several sister samples were produced, resulting in the availability of a "fresh" sample for each study), were exposed to 95% relative humidity at 60°C for 24 hours. With the use of a Nomarski microscope, no change in the Al₂O₃ sample was observed although the AlN showed signs of degradation.

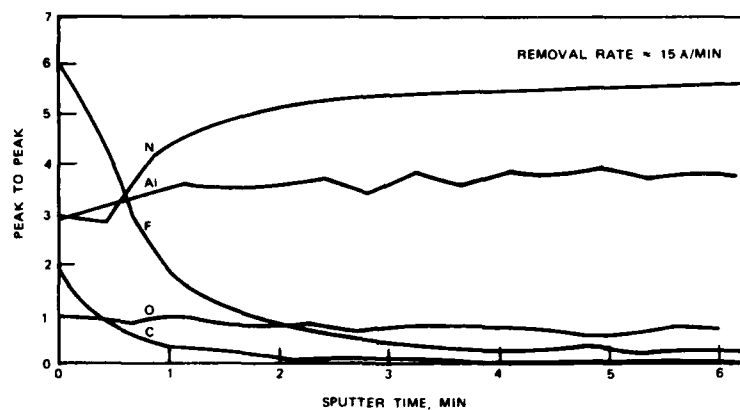


FIGURE 16. AES Profile of AlN Single Layer Exposed to Fluorine.
(a)

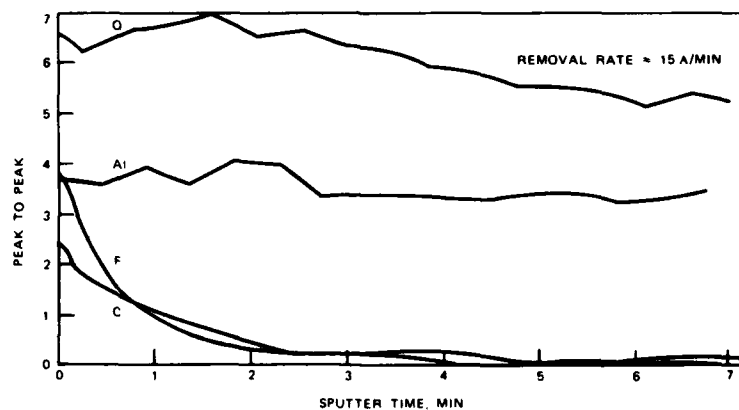


FIGURE 16. AES Profile of Al₂O₃ Single Layer Exposed to Fluorine.
(b)

CONCLUSIONS

The theoretical aspects of the AlN/Al₂O₃ multilayer designs indicate that perhaps the two-layer case may be more useful because of its bandwidth. At a reflectance loss of 0.5%, it has a bandwidth three times the size of that of the three-layer design. Although it has not been discussed in this study, simple calculations show that a two-layer quarter wave design is possible with a low reflectance (less than 0.005). This may be a more practical design. In this study, a film corresponding to the three-layer design was prepared and transmission results indicated a need to determine more accurately the indices of refraction and deposition rates of the materials.

Film structure studies indicate that defects are apparent on both the AlN, Al₂O₃, and AlN/Al₂O₃/AlN multilayer samples. The defects on the Al₂O₃ appear to be smaller (approximately 1 μ m in diameter as compared with 5-10 μ m for the AlN single and multilayer films). The growth structure of the AlN is columnar, whereas that of the Al₂O₃ is much smoother (no structure is observed at the SEM magnifications used in this study).

Laser damage studies indicate that multilayer damage initiates predominantly at defect sites. In addition, SAM studies indicate that the AlN outer layer of the multilayer oxidizes upon exposure to laser radiation in air. The amount of oxidation increases with increasing laser energy density and is also greater within damage craters than on portions of the film which appear to be unaffected by the laser radiation.

Fluorine exposure studies show depth penetrations of approximately 30-45 Å (assuming the same sputter removal rate of AlN and Al₂O₃ as SiO₂) in both films with the fluorine dropping off with depth at a slightly higher rate in the Al₂O₃.

Surface analysis techniques such as SEM and SAM indicate a sensitivity of the films to the electron beams. It is not certain whether the film changes after electron beam exposure are permanent.

After exposure to humidity, an Al₂O₃ single-layer film underwent no observable change, although a AlN single-layer film exposed to the same condition degraded.

Most importantly, it is apparent that the use of an AlN/Al₂O₃ multilayer in the environments described above dictates that Al₂O₃ should be the outer material.

Overall, both films appear to be durable and relatively easy to prepare as a multilayer system. These studies are continuing and clean room techniques are being applied to achieve films with fewer defects.

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APPENDIX

COMPUTER PROGRAM ANTIREF FOR CALCULATING
THREE-LAYER ANTIREFLECTION COATING SOLUTIONS

The following program calculates the thickness solutions for zero reflectance from a three-layer thin film system (based upon the work of Baer; reference 2).

```

C  ANTIREF.FOR
C
C  LESLIE G. KOSHIGOE
C
C  APRIL 9, 1984
C
C  THIS PROGRAM CALCULATES THE THICKNESS SOLUTIONS FOR ZERO
C  REFLECTANCE FROM A THREE LAYER THIN FILM SYSTEM BASED UPON
C  THE WORK OF A.D. BAER (REPORT ON MULTILAYER FILM PROGRAM FOR
C  JULY, 1976 - SEPT., 1977; NAVAL WEAPONS CENTER, CHINA LAKE,
C  CALIFORNIA). THE FILM IS ASSUMED TO BE LOSSLESS AND CERTAIN
C  INPUT PARAMETERS ARE REQUIRED (WAVELENGTH OF RADIATION AND
C  INDICES OF REFRACTION):
C
C  THE OPTICAL THICKNESSES FOR LAYERS ONE AND TWO (OUTER AND
C  MIDDLE LAYERS, RESPECTIVELY) ARE PLOTTED AGAINST THAT FOR THE
C  THIRD LAYER (ON SUBSTRATE). THUS, THE USER MAY CHOOSE A
C  DESIRABLE THIRD LAYER THICKNESS AND READ OFF THE THICKNESSES
C  NECESSARY FOR THE OTHER LAYERS. NOTE THAT THERE ARE TWO
C  SOLUTIONS FOR THE FIRST AND SECOND LAYER THICKNESSES FOR
C  EVERY THIRD LAYER THICKNESS (WHERE A SOLUTION IS POSSIBLE).
C  THESE SOLUTIONS ARE DISTINGUISHED BY SOLID AND DASHED LINES.
C  ALSO NOTE THAT SOLUTIONS EXIST FOR EVERY MULTIPLE OF A HALF-
C  WAVE FOR THE OPTICAL THICKNESSES.
C
C  THE PHYSICAL THICKNESSES ARE NOT CALCULATED IN THIS PROGRAM,
C  ALTHOUGH THEY MAY BE DETERMINED FROM:
C       $TI = (XI * WVLGTH) / NI$ 
C  WHERE TI DENOTES THE THICKNESS OF LAYER I, XI DENOTES THE
C  OPTICAL THICKNESS OF LAYER I, WVLGTH DENOTES THE WAVELENGTH
C  OF THE INCIDENT RADIATION, AND NI DENOTES THE INDEX OF REFRACTION OF LAYER I.

```

NWC TP 6651

```

C
C DEFINITION OF PARAMETERS:
C   WVLGTH = WAVELENGTH OF IMPINGENT RADIATION (FREE SPACE),
C           MICRONS,
C   N0 = INDEX OF REFRACTION OF ATMOSPHERE,
C   N1 = INDEX OF REFRACTION OF LAYER 1,
C   N2 = INDEX OF REFRACTION OF LAYER 2,
C   N3 = INDEX OF REFRACTION OF LAYER 3,
C   N4 = INDEX OF REFRACTION OF SUBSTRATE,
C   R10 = FRESNEL COEFFICIENT BETWEEN LAYER 1 AND ATMOSPHERE,
C   R21 = FRESNEL COEFFICIENT BETWEEN LAYERS 2 AND 1,
C   R23 = FRESNEL COEFFICIENT BETWEEN LAYERS 2 AND 3,
C   R34 = FRESNEL COEFFICIENT BETWEEN LAYERS 3 AND 4,
C   RA = FRESNEL COEFFICIENT BASED ON SMITH'S THEOREM,
C   RB = FRESNEL COEFFICIENT BASED ON SMITH'S THEOREM,
C   RBMAG = MAGNITUDE OF RB,
C   DELPA = PHASE OF RA (SOLID SOLUTION),
C   DELNA = PHASE OF RA (DASHED SOLUTION),
C   DELP = PHASE OF RB,
C   XT = VALUES OF THE THIRD LAYER THICKNESS FOR WHICH
C        CHECKS ON THE EXISTANCE OF SOLUTIONS ARE MADE,
C   X1POS = LAYER 1 OPTICAL THICKNESS (SOLID SOLUTION),
C   X1NEG = LAYER 1 OPTICAL THICKNESS (DASHED SOLUTION),
C   X2POS = LAYER 2 OPTICAL THICKNESS (SOLID SOLUTION),
C   X2NEG = LAYER 2 OPTICAL THICKNESS (DASHED SOLUTION),
C   X3 = LAYER 3 OPTICAL THICKNESS.
C
C DIMENSION X3(500),X1PCS(500),X1NEG(500),X2POS(500)
C DIMENSION X2NEG(500)
C REAL*4 N0,N1,N2,N3,N4
C CHARACTER*60 HEAD1
C
C READ IN INPUT PARAMETERS. THE WAVELENGTH (MICRONS) AND INDICES
C OF REFRACTION SHOULD BE STORED IN THE FILE "INDEX.DAT" PRIOR TO
C EXECUTION.
C
C OPEN(UNIT=10,FILE="ANTIREF.DAT",ACCESS="SEQUENTIAL",
C *   FORM="FORMATTED",STATUS="NEW")
C OPEN(UNIT=11,FILE="INDEX.DAT",ACCESS="SEQUENTIAL",
C *   FORM="FORMATTED",STATUS="OLD")
C
C READ(10) WVLGTH
C 10 FORMAT(F5.2)
C READ(10) N0
C READ(10) N1
C READ(10) N2
C READ(10) N3
C READ(10) N4
C 20 FORMAT(F6.4)
C
C CALCULATION OF FRESNEL COEFFICIENTS AT THE VARIOUS INTERFACES.
C
C R10=(N1-N0)/(N1+N0)
C R21=(N2-N1)/(N2+N1)
C R23=(N2-N3)/(N2+N3)
C R34=(N3-N4)/(N3+N4)
C
C PI=3.1415926535

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NWC TP 6651

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C      CALCULATION OF THE SOLUTIONS (OR OF THE X1 AND X2 VALUES)
C      FOR EACH X3 VALUE.
C
      J=1
      XT=0.
      DO 40 I=1,501
      REMAG=(SQRT((R23+R24**2**2**2**2*(COS(4.*PI*XT))+R34**2*(COS(4.*PI*XT)
**R23**R34**2)**2**2*(SIN(4.*PI*XT))**2*(R23**2-1.)**2))/
*(1.+2.*R23**R24**2*(COS(4.*PI*XT))+R23**2**R34**2)
C
      FUNC=(R21**2+R10**2-REMAG**2*(1.+R21**2**R10**2))/(2.*R21**R10*
*(REMAG**2-1.))
C
C      DETERMINATION OF WHETHER A SOLUTION EXISTS FOR THE PRESENT XT
C      VALUE. IF THERE IS NO SOLUTION, THEN XT IS INCREMENTED UP BY
C      0.001.
C
      IF(ABS(FUNC).GT.1.) THEN
      GO TO 30
      END IF
C
C      A SOLUTION EXISTS SO X3 IS SET EQUAL TO XT, AND THE X1 AND X2
C      VALUES ARE CALCULATED.
C
      X3(J)=XT
      X1POS(J)=(1./(4.*PI))*ACOS(FUNC)
      X1NEG(J)=-X1POS(J)
C
C      DETERMINES THE QUADRENT THAT PHASE ANGLE DELPA IS IN.
C
      DELNUM=(-R10*STN(4.*PI*X1POS(J)))*(1.-R21**2)
      DELDEN=R21+R21**2**R10*(COS(4.*PI*X1POS(J)))+R10
**COS(4.*PI*X1POS(J))+R21**R10**2
      IF(DELDEN.EQ.0.) GO TO 900
      DELPA=ATAN(ABS(DELNUM/DELDEN))
      IF(DELNUM.GT.0.) THEN
      IF(DELDEN.GT.0.) GO TO 2000
      END IF
      IF(DELNUM.GT.0.) THEN
      IF(DELDEN.LT.0.) GO TO 1000
      END IF
      IF(DELNUM.LT.0.) THEN
      IF(DELDEN.LT.0.) GO TO 1100
      END IF
      IF(DELNUM.LT.0.) THEN
      IF(DELDEN.GT.0.) GO TO 1200
      END IF
      IF(DELNUM.EQ.0.) THEN
      IF(DELDEN.GT.0.) GO TO 1300
      ELSE
      IF(DELDEN.LT.0.) GO TO 1400
      END IF
      900 IF(DELDEN.EQ.0.) THEN
      IF(DELNUM.GT.0.) GO TO 1500
      ELSE
      IF(DELNUM.LT.0.) GO TO 1600
      END IF
      1000 DELPA=PI-DELP
      GO TO 2000
      1100 DELPA=PI+DELP
      GO TO 2000
      1200 DELPA=2.*PI-DELP
      GO TO 2000
      1300 DELPA=0.
      GO TO 2000
      1400 DELPA=PI
      GO TO 2000
      1500 DELPA=PI/2.
      GO TO 2000
      1600 DELPA=(3.*PI)/2.

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NWC TP 6651

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C
C      DETERMINES THE QUADRENT THAT PHASE ANGLE DELB IS IN.
C
2000 DELNUM=(-R34*SIN(4.*PI*X3(J)))*(1.-R23**2)
    DELDEN=(R23+R23**2*R34*(COS(4.*PI*X3(J)))+R34*COS(4.*PI*
    *X3(J))+R23*R34**2)
    IF(DELDEN.EQ.0.) GO TO 2900
    DELB=ATAN(ABS(DELNUM/DELDEN))
    IF(DELNUM.GT.0.) THEN
        IF(DELDEN.GT.0.) GO TO 4000
    END IF
    IF(DELNUM.GT.0.) THEN
        IF(DELDEN.LT.0.) GO TO 3000
    END IF
    IF(DELNUM.LT.0.) THEN
        IF(DELDEN.LT.0.) GO TO 3100
    END IF
    IF(DELNUM.LT.0.) THEN
        IF(DELDEN.GT.0.) GO TO 3200
    END IF
    IF(DELNUM.EQ.0.) THEN
        IF(DELDEN.GT.0.) GO TO 3300
    ELSE
        IF(DELDEN.LT.0.) GO TO 3400
    END IF
2900 IF(DELDEN.EQ.0.) THEN
    IF(DELNUM.GT.0.) GO TO 3500
ELSE
    IF(DELNUM.LT.0.) GO TO 3600
END IF
3000 DELB=PI-DELB
    GO TO 4000
3100 DELB=PI+DELB
    GO TO 4000
3200 DELB=2.*PI-DELB
    GO TO 4000
3300 DELB=0.
    GO TO 4000
3400 DELB=PI
    GO TO 4000
3500 DELB=PI/2.
    GO TO 4000
3600 DELB=(3.*PI)/2.
4000 X2PCS(J)=(1./(4.*PI))*(DELF+DELB)
C
C      DETERMINES THE QUADRENT THAT PHASE ANGLE DELNA IS IN.
C
    DELNUM=(-R10*SIN(4.*PI*X1NEG(J)))*(1.-R21**2)
    DELDEN=(R21+R21**2*R10*(COS(4.*PI*X1NEG(J)))+R10*COS(4.*PI*
    *X1NEG(J))+R21*R10**2)
    IF(DELDEN.EQ.0.) GO TO 4900
    DELNA=ATAN(ABS(DELNUM/DELDEN))
    IF(DELNUM.GT.0.) THEN
        IF(DELDEN.GT.0.) GO TO 6000
    END IF
    IF(DELNUM.GT.0.) THEN
        IF(DELDEN.LT.0.) GO TO 5000
    END IF
    IF(DELNUM.LT.0.) THEN
        IF(DELDEN.LT.0.) GO TO 5100
    END IF
    IF(DELNUM.LT.0.) THEN
        IF(DELDEN.GT.0.) GO TO 5200
    END IF
    IF(DELNUM.EQ.0.) THEN
        IF(DELDEN.GT.0.) GO TO 5300
    ELSE
        IF(DELDEN.LT.0.) GO TO 5400
    END IF

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NWC TP 6651

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4900 IF(DEL DEN.EC.0.) THEN
      IF(DEL NUP.GT.0.) GO TO 5500
      ELSE
        IF(DEL NUP.LT.0.) GO TO 5600
      END IF
5000 DELNA=PI-DELNA
      GO TO 6000
5100 DELNA=PI+DELNA
      GO TO 6000
5200 DELNA=2.*PI-DELNA
      GO TO 6000
5300 DELNA=0.
      GO TO 6000
5400 DELNA=PI
      GO TO 6000
5500 DELNA=PI/2.
      GO TO 6000
5600 DELNA=(3.*PI)/2.
6000 X2NEG(J)=(1./(4.*PI))*(DELNA+DEL2)
      J=J+1
      30 XT=XT+0.001
      40 CONTINUE
C
      DO 42 I=1,J-1
        X1NEG(I)=X1NEG(I)+0.5
        X2POS(I)=X2POS(I)-0.5
      42 CONTINUE
C
      THE OPTICAL THICKNESSES ARE WRITTEN TO THE FILE "ANTIREF.DAT".
C
      WRITE(10,43)WVLGTH,N0,N1,N2,N3,N4
43  FORMAT(1X,"WAVELENGTH = ",F5.2,"MICRONS",/,1X,"N0 = ",F6.4,/,
        *1X,"N1 = ",F6.4,/,1X,"N2 = ",F6.4,/,1X,"N3 = ",F6.4,/,1X,"N4 = ",
        *F6.4,///)
      WRITE(10,45)
45  FORMAT(3X,"X3",5X,"X1POS",3X,"X1NEG",3X,"X2POS",3X,"X2NEG",/)
      WRITE(10,50)(X3(I),X1POS(I),X1NEG(I),X2POS(I),X2NEG(I),I=1,J-1)
50  FORMAT(1X,F7.4,1X,F7.4,1X,F7.4,1X,F7.4,1X,F7.4)
C
      PLOTTING SECTION. (USES DISPLA GRAPHICS.)
C
      CALL COMPRS
      CALL PAGE(9.5,11.0)
      CALL SWISSM
      CALL SHDCMP(90.,1,0.01,1)
      CALL AREA2D(5.,5.)
      CALL FRAME
C
      WRITES AXIS LABELS AND THE PLOT HEADING.
C
      CALL RESET("MIXALF")
      CALL XNAME("X3 ((Wavelengths))",100)
      CALL YNAME("X1 and X2 ((Wavelengths))",100)
      CALL MIXALF("L/CGREEK")
      WRITE (HEAD1, 40) WVLGTH
60  FORMAT("Antireflectance Coating ((", F4.2, "(M)m)")
      CALL HEADIN(XREF(HEAD1), 100, 1.5, 1)

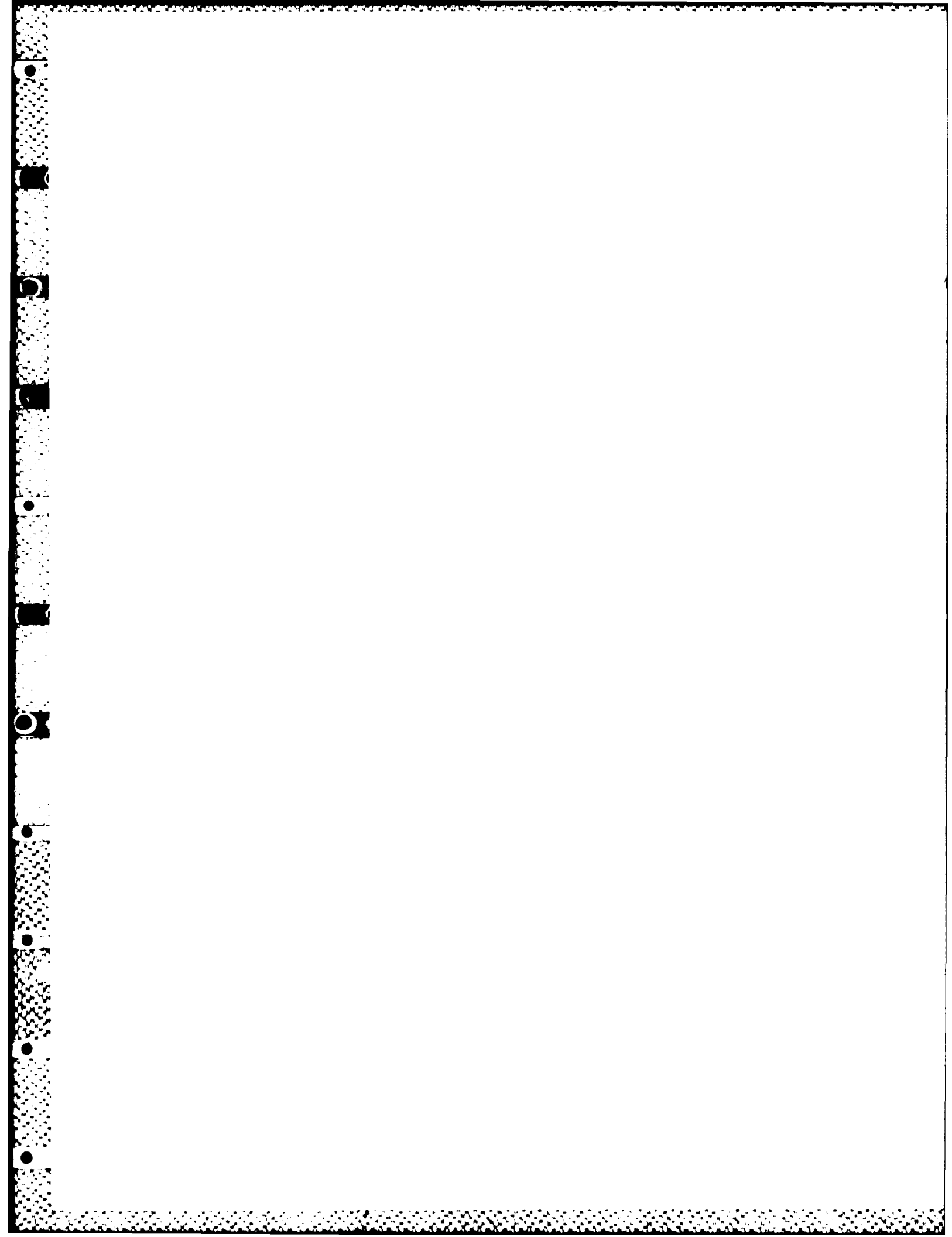
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NWC TP 6651

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C
C      PLOTS X1 AND X2 VS. X3.  A SPLINE CURVE FITTING ROUTINE IS USED
C      FOR FITTING THE DATA.
C
CALL GCAF(0.,0.1,0.5,0.,0.1,0.5)
CALL RASPLN(0)
CALL T-MERV(0.03)
CALL B-(JRV(.02,.02))
CALL CURVE(X3,X1PCS,J-1,C)
CALL CURVE(X3,X2PCS,J-1,C)
CALL DASH
CALL CURVE(X3,X1NEG,J-1,0)
CALL CURVE(X3,X2NEG,J-1,0)
CALL RESET("DASH")
CALL ENDCPL(0)
CALL DONEPL
END
```

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